

# Domains Formation Mediated by Electromagnetic Fields in Very Dilute Aqueous Solutions: 1. Quantum Electrodynamical Aspects

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For a list of abbreviations, see **Table 1** on page 45.

(This is the first part of a three-part series.)

## Abstract

$10^{-7}$  -  $10^{-4}$  m sized molecular groupings in solutions are the foci of our research. The concentration of the solutions varies from 2 M to  $10^{-20}$  M. The solutions are prepared by serial dilutions of a stock solution with polar solvents (*e.g.*, water) and vigorous shaking after each dilution step. Light scattering, dielectric permittivity, electric conductivity, pH and others measurements, reported in previous publications, evidenced: (i) electromagnetic fields mediate formation of the groupings; (ii) the impact of the groupings on some of the solutions' physicochemical properties (*e.g.*, electric conductivity) is correlated with their bio-activity. The aim of our current research is analyzing experimental data pertaining to the groupings' characteristics and confirming that these agree with those predicted by quantum electrodynamics (QED). Towards this aim, in the current paper we provide a concise overview of recently derived predictions of QED relevant to molecular

groupings in solutions. We also cite publications presenting experimental data verifying some of these predictions. In the two following papers in this journal's issue, we employ QED for analyzing recent published experimental data pertaining to groupings in solutions of, respectively, strong electrolytes and other solutes (weak electrolytes, non-electrolytes).

## Introduction

Down to ultra low concentrations ( $C$ ) - about  $10^{-20}$  M - serial diluted vigorously shaken polar liquids (SDVSPL) may affect bio-systems (Palmina *et al.*, 1994; Konovalov *et al.*, 2008, 2014a; Konovalov and Ryzhkina 2014b).<sup>a</sup> The biological active solutes in-

<sup>a</sup> SDVSPL preparation involves serial decimal or centesimal diluting a "stock" solution. Hitherto, mainly aqueous SDVSPL have been studied, but solutions with other solvents were also investigated. Aqueous SDVSPL are prepared with freshly doubly distilled water or water purified by Simplicity® Water Purification Systems - Millipore, with specific electrical conductivity below 2.5  $\mu\text{S}/\text{cm}$ . Dust

clude inorganic-, organic- or bio-molecules, strong or weak electrolytes. For example, pesticides, poisons, synthetic or natural remedies affecting enzymes, membranes and organisms. In ultra low concentrated (ULC) SDVSPL at ambient conditions, part of the solvent molecules form groupings, as initially revealed with calorimetric, electric conductivity and pH data; serial diluted solutions, which at each dilution step are not vigorously shaken, do not contain groupings for  $C$  below a solute type dependent critical concentration ( $C_{crit}$ ); typically  $\sim 10^{-8} \text{ M} < C_{crit} < \sim 10^{-6} \text{ M}$  (Elia and Niccoli, 1999, 2000, 2004a, 2004b). The groupings comprising  $10^{-7} - 10^{-5} \text{ m}$  sized molecular associates, their electrokinetic potential and their impact on the liquid's dielectric permittivity were first uncovered with dynamic light scattering, electrophoresis and dielcometric titrations, respectively (Konovalov *et al.*, 2008; Ryzhkina *et al.*, 2009). Stabilization of the groupings occurs during the period of about 1-18 hours after the SDVSPL's preparation. For  $10^{-13} \text{ M} < C < 10^{-3} \text{ M}$  aqueous SDVSPL of strong electrolytes, transmission electron microscopy, atomic force microscopy and dielectric permittivity data revealed presence of  $10^{-5} - 10^{-4} \text{ m}$  sized molecular associates (Lo *et al.*, 1996, 2009; Ho, 2014).

For solute type dependent  $C$  ranges, correlation was observed between the solutions' bioactivity, electric conductivity, the groupings' effective hydrodynamic diameter and

is removed. Stock solutions are analyzed for absence of impurities.  $C$  of stock solutions are in the  $4 \text{ M} - 10^{-3} \text{ M}$  range. After each dilution step, SDVSPL are vigorous shaken, *e.g.*, with lab dancer shaker, by vertical vortexing or other methods. Plastic or glass vessels are used. Temperature and pressure are kept constant, typically, respectively, at  $298 \text{ K}$  and  $\sim 1 \text{ Atm}$ . As controls, the solvent (*e.g.*, doubly distilled water) is serial diluted and shaken after each dilution step, with all experimental parameters identical to those of SDVSPL preparation.

electrokinetic potential; all these properties non-linearly depend on  $C$  and are reproducible (Konovalov *et al.*, 2014c). Not for all kinds of solutes, their ULC SDVSPL contain groupings. For some solute types their ULC SDVSPL have the customary characteristics of infinite diluted solutions (Konovalov, 2013). Solute attributes required for inducing groupings and the origins of their physical, chemical, catalytic and bio-active properties are not yet clarified. Impurities released by containers affect SDVSPL but cannot account for their typical properties. For a concise discussion of impurities' effects see Yinnon and Liu (2015a).

As to forces underlying the groupings formation, electrodynamic ones play a role -- electromagnetic fields (EMF) affect aqueous SDVSPL containing groupings (Montagnier *et al.*, 2009; Ryzhkina *et al.*, 2011; Elia *et al.*, 2012; Konovalov *et al.* 2014a,b). On storing samples under hypo electromagnetic conditions, *i.e.*, in a Permalloy container with residual field of 10 nano Tesla, no groupings are observable for  $10^{-20} \text{ M} < C < C_{thr}$ ; weak EMF influence the groupings for  $C_{thr} < C < \sim 10^{-4} \text{ M}$ . The threshold concentration  $C_{thr}$  is solute type dependent, typically  $10^{-10} \text{ M} < C_{thr} < 10^{-6} \text{ M}$ .

Customary models of ULC solutions at ambient conditions cannot account for SDVSPL's properties. These models predict that: EMF, serial dilutions or vigorous shaking do not affect SDVSPL characteristics; solvated solutes distribute homogeneously, move independently and randomly; polar solvent molecules (except solvation shells' solvent molecules) move randomly; for aqueous solutions, its water molecules ( $\text{H}_2\text{O}$ ) form flickering hydrogen-bond networks (Horne, 1971). These customary models explicitly include electrostatic forces and *assume electrodynamic ones can be treated perturbatively*. However, quantum electrodynamic (QED) models *explicitly* including electrodynamic forces show inter-

actions between EMF and liquid molecules may lead to formation of various domains (Del Giudice, 1988, 2000; Arani *et al.*, 1995; Preparata, 1995 chapters 2, 5, 10; Yinnon and Yinnon, 2012). For example, EMF interactions with electrolytic solutes, polar solute molecules or with solvent molecules with sufficiently large electric dipole moments, for solute type dependent  $C$  ranges, may lead to distinctive domain types. Within the context of formal QED theory of polar liquids, the conditions for formation of these domains and their properties were *ab initio* derived. SDVSPL observed phenomena signify QED has to be employed for their explanation. Indeed, the QED model for SDVSPL proposed by Yinnon and Yinnon (2011) has provided consistent explanations for various phenomena, *e.g.*, self-organization of molecules in SDVSPL, these liquids' electric conductivity, heat of mixing and their dependence on time and volume (Yinnon and Elia 2013).

Our study's foci are: (a) employing QED for explaining recently observed (and to the best of our knowledge yet unexplained) characteristics of the various groupings present in SDVSPL; (b) elucidating the relation between these characteristics and the liquid's properties, *e.g.*, its dielectric permittivity, spectra and bioactivity. As to their importance, SDVSPL have implications for numerous technologies and bio-processes.

Since QED of polar liquids hitherto mainly is employed for explaining special phenomena, many readers may be unfamiliar with it. Its aspects relevant to our analyses of SDVSPL were derived in a series of publications published since 1988. Some of these publications contain many mathematical expressions. Therefore we deem it appropriate to provide a short overview of the topic to the non-physicists, *i.e.*, an overview which uses intuitive arguments only. Of course avoidance of the rigorous mathematical technicalities implies we have

to pay the price of some imprecision. For those interested to study the issues more in depth, we point out their sources.

Based on the QED properties presented in our overview, Yinnon and Yinnon (2011) and Yinnon and Elia (2013) proposed a model for very dilute solutions. We will employ their model for explaining recently measured SDVSPL properties. Including in this paper a concise summary of their model and employing it for analyzing measured SDVSPL properties would make it excessively lengthy. Therefore, below we suffice with providing an overview of QED properties of polar liquids. In the two publications following this one (Yinnon and Liu, 2015a, 2015b) respectively, for SDVSPL of strong electrolytes and of other solutes (weak electrolytes, non-electrolytes), we will present their models and employ these for explaining recently observed phenomena.

As to the outline of this paper, we first present an overview of polar liquids' characteristics predicted by QED. The overview starts with some historical milestones pertaining to modeling effects of EMF on intermolecular interactions and thermodynamics of molecular ensembles. Next it focuses on structural aspects of the ensembles resulting from electrodynamic interactions, *i.e.*: phase transitions leading to formation of coherence domains; the physics underlying these domains; the domains' effects on thermodynamics of the ensembles; formation of supra-domains; types of domains present in polar liquids or their solutions; properties of the various domains hitherto identified. After that, in the Discussion section, we shortly discuss the status of experimental data confirming characteristics of polar liquids predicted by QED. In the Conclusion section we point out some challenges ahead. A list with abbreviations we add at the end of this paper. We stress that in the following we do not derive any new features of QED.

## QED Predictions for Polar Liquids' Characteristics

### *Modeling Effects of EMF on Intermolecular Interactions*

Assessing effects of electrodynamic forces on thermodynamic and structural properties of molecular ensembles for many decades posed numerous challenges. London (1930) showed that quantum fluctuations of the EMF affect intermolecular forces. Such fluctuations can be represented by virtual photons popping out of empty space (the quantum vacuum), as allowed by the uncertainty principle. These fluctuations are non-thermal -- these persist when temperature vanishes. London applied second order perturbation techniques to the electrostatic interactions between molecular dipoles to account for quantum fluctuations. His approach was a non-relativistic one and pertinent only to interactions between molecules in vacuum over distances of the order of  $10^{-10}$ - $10^{-9}$  m, *e.g.*, rarefied gas. He derived that the fluctuations underlie the intermolecular dispersion force, which is a component of the van der Waals force. Casimir and Polder (1948) used fourth-order perturbation techniques to derive the interactions between molecules and the EMF over longer distances, *i.e.*, distances of the order or larger than the molecules' characteristic absorption wavelengths. At such distances, relativistic effects play a role. Their inclusion for example enabled explaining the Casimir effect, which exemplifies the considerably strength of the interactions mediated by virtual photons (Casimir 1948). After the advent of quantum field theory, which proved to be the most accurate physics theory, Dzyaloshinskii *et al.* (1961) employed it to study the interactions between molecules in a liquid and the EMF. Their approach is based on perturbation techniques and Feynman diagrams. It for example enabled them to derive the effects of dispersion forces on the thermodynamics of liquids.

### *Phase Transitions Induced by QED Interactions Leading to Coherence Domains*

The aforementioned approaches treated interactions between molecules and the EMF as small perturbations. Albeit conditions for fluctuating EMF majorly affecting condensed matter, *e.g.*, leading to phase transitions, were predicted (Dicke, 1953, Hepp and Lieb, 1973a,b). However, issues like the Hamiltonian's form required for describing the interactions between the fluctuating EMF and molecules (Bialynicki-Birula and Rzazewski, 1979) for many years prevented their verification. Hence, a practical approach was adopted: condensed matter was modeled by explicitly including electrostatic interactions, while QED ones were described perturbatively or even ignored. The resulting theories explained numerous phenomena and led to the current customary models of polar liquids. Yet Preparata (1988, 1995) and Del Giudice *et al.* (1988, 1993) succeeded to resolve the aforementioned issues. They developed a non-perturbative (variational) quantum field theory model of liquids using the Feynman path integral. Their model enabled investigating phase transitions. They derived that QED interactions may lead to phase transitions whenever density and temperature are, respectively, above and below transition values. These values depend on variables like the coupling between matter and radiation. For water and other polar liquids, they showed that the transitions lead to auto-organization of a fraction of the liquid's molecules in  $10^{-7}$ - $10^{-4}$  m sized coherence domains (CD). This fraction depends on temperature. Within a CD, solvent or solute molecules coherently oscillate in-phase with a coherently condensed EMF. As to the molecules which do not join CD, these move randomly in the interstices between the CD. Independent research groups confirmed aforementioned (Sivasubramanian *et al.*, 2001, 2002, 2003, 2005; Emary and Brandes, 2003; Apostel, 2009)



### *The Physics Underlying Coherence Domains*

For elucidating the physics underlying CD, we summarize explanations by Del Giudice *et al.* (1998), Preparata (1995 chapters 2,3,10), Arani *et al.* (1995) and Bono *et al.* (2012). Consider an ensemble of  $N$  molecules, which initially all move randomly within a volume  $V_e$ . Whenever the energy of a quantum fluctuation of the EMF (a virtual photon) equals that of a specific molecular excitation, all molecules within the volume spanned by the photon are candidates for excitation. This volume we denote  $V_{\text{photon}}$ ; it equals  $\lambda^3$ , with  $\lambda$  denoting the photon's wavelength.  $V_{\text{photon}}$  contains about  $10^6 - 10^{18}$  molecules, because electronic, vibrational or rotational excitations of molecules require  $10^{-7} < \lambda < 10^{-4}$  m, while molecular sizes typically are about  $10^{-9} - 10^{-10}$  m. The photon excites one molecule with probability  $P_r$ . The decay time of the excited state determines the duration of the molecule's excitation. On de-excitation, the photon emitted can fly away or excite a second molecule. The probability  $P_N$  that it excites another molecule is  $P_N = P_r \lambda^3 (N / V_e)$ . For  $P_N < 1$ , the photon eventually returns to the quantum vacuum. However, when  $P_N = 1$  [a condition occurring when  $N / V_e$  is larger than a transition density  $(N / V_e)_{\text{trans}}$ ], the photon loses the chance of leaving  $V_{\text{photon}}$  and bounces from one molecule to another. Additional virtual photons may share the same fate, leading to a sizable EMF condensing within  $V_{\text{photon}}$ .<sup>b</sup> The condensed EMF causes all

<sup>b</sup> The condensation results from renormalization of the photon's oscillation time when it interacts with the molecules. Due to the decay time of the molecules' excited state, the frequency of a photon bouncing between molecules in a CD is smaller than the frequency of the free photon. According to the Einstein equation, the differences in these frequencies causes the photons' mass (which is zero for the free photon) to become imaginary. That is, the photons are unable to propagate, these are trapped as excitations of the molecules.

molecules within  $V_{\text{photon}}$  to oscillate coherently between their ground state  $|0\rangle$  and an excited state  $|b\rangle$ . Accordingly,  $V_{\text{photon}}$  was denoted Coherence Domain (CD). Since  $V_e$  can be partitioned into many  $V_{\text{photon}}$ , other virtual photons may lead to more CD formations. Virtual photons with different  $\lambda$  initially may excite molecules to different  $|b\rangle$ . However the system's dynamics ultimately cause the molecules in all CD to oscillate between  $|0\rangle$  and the same  $|b\rangle$ . The coupling between the molecules and the EMF determines  $|b\rangle$ .

### *Thermodynamics of Molecular Ensembles Containing Coherence Domains*

Only in open systems do CD form (Preparata, 1995 chapters 2,3,10; Arani *et al.*, 1995). The energy of molecules in CD is lower than that of the ensemble's randomly moving molecules. Accordingly, the energy gained on inclusion of a molecule within a CD has to be released to the environment as the phase transition's latent heat. For ensembles in contact with a heat bath, temperature ( $T$ ) determines the fraction of molecules included in CD. This fraction is an inverse function of  $T$ . Only below a transition temperature CD form. The energy a molecule gains on its inclusion within a CD is an inverse function of its distance from the center of its CD. When this energy is only slightly less than  $k_B T$  (with  $k_B$  the Boltzmann constant), the disruptive dynamics of thermal collisions may push molecules out of the CD. Therefore in specific temperature ranges, CD are metastable with molecules constantly desorbing and adsorbing at its surface.

### *Supra-coherence Domains*

CD may agglomerate into supra-domains (supra-CD) (Preparata, 1995, chapter 10). Supra-CD are not ensembles of molecules but agglomerates of domains, like domains in liquid crystals. The agglomeration of CD is energetically favorable, because close packing of CD enables the evanescent tails

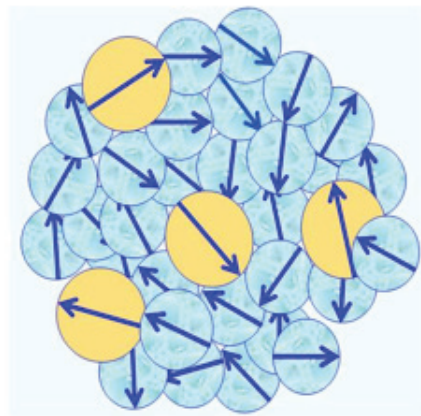
of their condensed EMF to overlap and interfere constructively. As a result, within a supra-CD the oscillations of all molecules are coherent.

### *Types of Coherence Domains Present in Polar Liquids or Their Solutions*

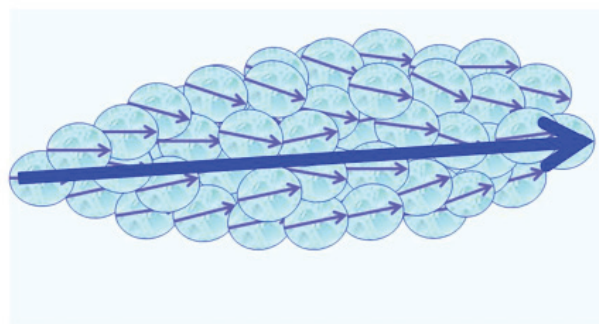
EMF interactions with electrolytic solutes, polar solute molecules or with solvent molecules with sufficiently large electric dipole moments, for solute type dependent  $C$  ranges, may lead to various CD types (Del Giudice, 1988, 1998, 2000; Arani *et al.*, 1995; Preparata, 1995 chapters 2, 5, 10; Yinon and Yinon, 2012). The CD types and those of their properties required for modeling SDVSPL we summarize in the next paragraphs, using nomenclature of our previous publications. Hitherto identified CD include:

- $CD_{elec}$ , which are composed of solvent molecules only. Some of these molecules reside in an electronically excited state (see Figure 1). Their diameter is  $\sim 10^{-7}$  m.
- $CD_{rot}$ , which are composed of ferroelectric ordered polar solvent molecules (see Figure 2). Their diameter is of the order of  $\sim 10^{-4} - 10^{-5}$  m.
- $CD_{plasma}$ , which are composed of few solvated solutes and numerous polar solvent molecules (see Figure 3a). Their diameter is  $\sim 10^{-6}$  m. The plasma oscillations of their solute molecules are coherent.
- $IPD_{plasma}$ , which are composed of few solvated solutes and numerous polar solvent molecules (see Figure 3b). Their diameter is  $\sim 10^{-6}$  m. The plasma oscillations of the solute molecules within  $IPD_{plasma}$  are in phase, *i.e.*, an  $IPD_{plasma}$  is a special CD -- an In-Phase Domain.

To give the reader an intuitive feeling for the relative sizes of the  $CD_{rot}$ ,  $CD_{plasma}$  and  $CD_{elec}$ , we note their ratios are similar, respectively, to those of the sun, earth and moon.

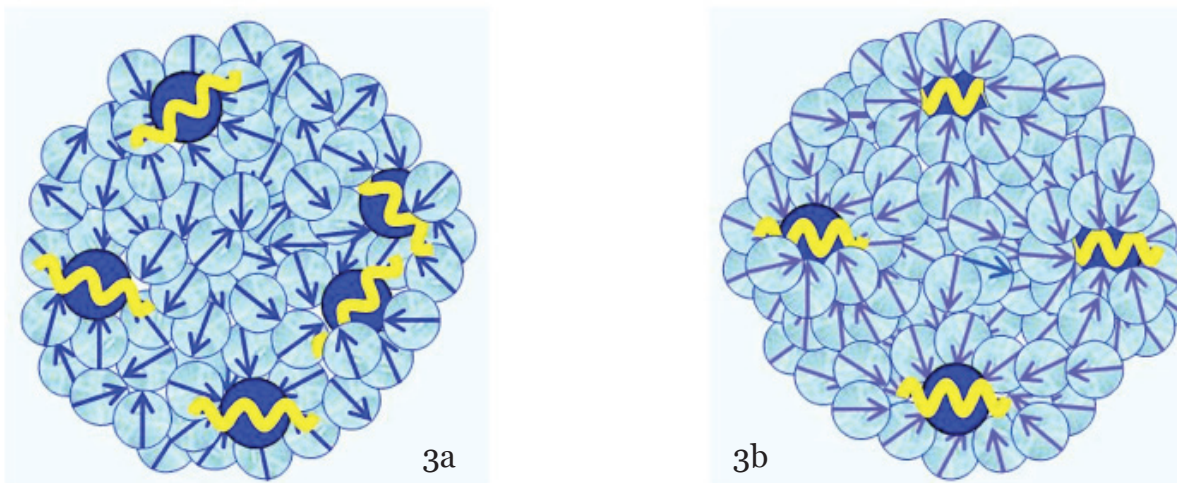


**Figure 1:** Schematic view of a  $CD_{elec}$  and its internal structure. Light-blue and yellow colored balls symbolize its molecules residing, respectively, in their ground and an excited electronic states. The molecules' electric dipole moments, symbolized by arrows, are randomly oriented.



**Figure 2:** Schematic view of a  $CD_{rot}$  and its internal structure. Light-blue colored balls symbolize the molecules constituting the domain. Their blue arrows symbolize the molecules' electric dipole moments. The large blue arrow symbolizes the electric dipole moment of the domain, resulting from the ferroelectric ordering of its molecules.

**Properties of  $CD_{elec}$**  --  $CD_{elec}$  formation is mediated by ultra violet EMF (Arani *et al.*, 1995; Preparata, 1995 ch. 10, Bono *et al.*, 2012). Their formation is a central aspect of the condensation of a liquid from its vapor. Only for water the characteristics of  $CD_{elec}$  have been derived. We denote with  $CD_{elec}^{H_2O}$  the  $CD_{elec}$  of water. The  $H_2O$  constituting  $CD_{elec}^{H_2O}$  reside in a state which is a superposition of their ground  $|0\rangle$  state (with a weight of about 87%) and the excited  $|b\rangle$  state (with a weight of about 13%). One electron of an



**Figure 3 a-b:** Schematic pictures of, respectively,  $CD_{plasma}$ , and  $IPD_{plasma}$  and their internal structures. Blue balls symbolizes solute molecules, with their yellow cosine curves symbolizing their plasma oscillations. For  $IPD_{plasma}$  these oscillations are in-phase, for  $CD_{plasma}$  these are just coherent. Light-blue colored balls symbolize polar solvent molecules. Their blue arrows symbolize their electric dipole moments. In  $IPD_{plasma}$  these dipoles are spherically symmetric oriented around their nearest neighbor solute molecule. In  $CD_{plasma}$  only the dipoles of  $H_2O$  constituting the hydration shell are spherically symmetric oriented; all non-hydration shell solvent molecules are oriented randomly.

$H_2O$  residing in its  $|b\rangle$  state is almost free (binding energy  $\approx 0.4$  eV). Hence, a  $CD_{elec}^{H_2O}$  is a pool of  $\sim 10^6$  quasi-free electrons located at their boundary with  $\sim 7, 21, 35, 49, \dots$  kHz excited states, and correspondingly an ensemble of quasi free protons (the partners of the quasi-free electrons (Del Giudice *et al.*, 1998, 2007)).  $H_2O$  are tetrahedrally ordered in  $CD_{elec}^{H_2O}$ .  $CD_{elec}^{H_2O}$  agglomeration in supra- $CD_{elec}^{H_2O}$  underlie the hydrogen-bond network of bulk water.

Energy gained by a  $H_2O$ 's inclusion in  $CD_{elec}^{H_2O}$  is  $\sim 0.17$  eV at  $T=273$  K and pressure of 1 Atm.  $CD_{elec}^{H_2O}$  only exist for temperatures below the transition temperature of 500 K. The fraction of  $H_2O$  included in  $CD_{elec}^{H_2O}$  for  $T < 180$  K equals 1. For  $180$  K  $< T < 500$  K, this fraction is less than 1. The  $H_2O$ , which are not included in  $CD_{elec}^{H_2O}$ , move randomly in between or evaporate. Chemical potentials determine the fraction of  $H_2O$  included within  $CD_{elec}^{H_2O}$ , percentages of  $H_2O$  randomly moving in between  $CD_{elec}^{H_2O}$  and percentages of evaporated  $H_2O$ . Bulk water and ice- $I_h$  contain  $CD_{elec}^{H_2O}$ . At  $T=298$  K, the fraction of  $H_2O$  included in  $CD_{elec}^{H_2O}$  is about 20 percent.

In bulk water, at ambient conditions, some  $H_2O$  continually adsorb on  $CD_{elec}^{H_2O}$  while simultaneously others desorb from the domains, causing a  $\sim 10^{-14}$  s timescale flickering landscape. Thus  $CD_{elec}^{H_2O}$  observation requires fast resolution probes. Interfaces,  $CD_{rot}$ ,  $CD_{plasma}$  and  $IPD_{plasma}$  may stabilize  $CD_{elec}^{H_2O}$  and supra- $CD_{elec}^{H_2O}$ , *i.e.*, reduce their flickering, easing their observation.  $CD_{elec}^{H_2O}$  cannot contain solutes. Solvated solutes locate outside  $CD_{elec}^{H_2O}$ .

**Properties of  $CD_{rot}$**  --  $CD_{rot}$  formation results from their solvent molecules' dipole moment interacting with Far Infra Red EMF (Del Giudice *et al.*, 1988; Del Giudice and Vitiello, 2006).  $CD_{rot}$  have an electric dipole moment due to the ferroelectric ordering of their solvent molecules (see Figure 2). In bulk water, as well as in the bulk of many other polar liquids, at ambient conditions  $CD_{rot}$  do not auto-organize. However, immersing objects with sizable asymmetric charge distributions (*e.g.*, macromolecules, hydrophilic membranes) may induce their formation, resulting in a permanent time dependent polarization. Solutes are pulled

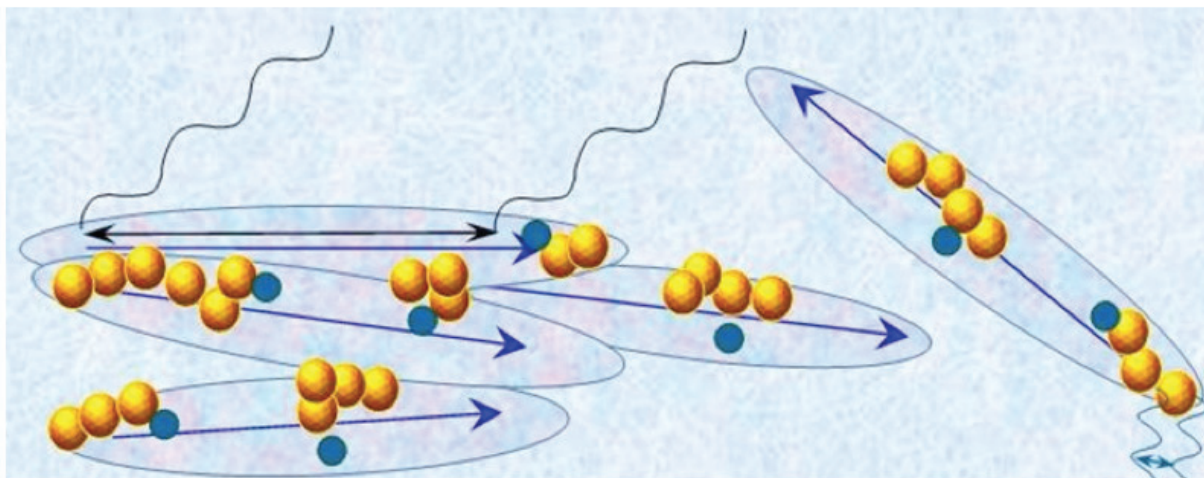


into  $CD_{rot}$ . Few solute particles can locate in  $CD_{rot}$  and do not wreck their host. Many solute molecules destroy  $CD_{rot}$ . Solute type determines critical  $C$  below which  $CD_{rot}$  persists ( $C_{crit}^{CD_{rot}}$ ).  $CD_{rot}$ 's molecules coherently oscillate between two rotational states, resonating in-phase with the coherently condensed photons mediating their interactions.

QED  $CD_{rot}$  are analogues of classical ferroelectric domains, which form due to an instability in the expression relating polarization density, polarization susceptibility and an external electric field (Sivasubramanian *et al.*, 2005). QED, classical electrodynamic models, computer simulations and experimental data show that for bulk polar liquids at pressure of 1 Atm,  $CD_{rot}$  may exist for temperatures below a transition temperature (Yinnon *et al.*, 2015c -- submitted to Water journal). For  $H_2O$  the value of this transition temperature is 1160 K. Whenever the energy a solvent molecule gains by

inclusion in  $CD_{rot}$  is higher than  $k_B T$ , the solvent molecules auto-organize into  $CD_{rot}$ . Whenever this energy is lower than  $k_B T$ ,  $CD_{rot}$  do not form spontaneously. However, by disturbing the liquid, for example by immersing a hydrophylic membrane into it,  $CD_{rot}$  may form. On removal of the membrane, these  $CD_{rot}$  are meta-stable and may persist for many months.

$CD_{rot}$  may agglomerate into supra-domains. Supra- $CD_{rot}$  may stabilize supra- $CD_{elec}^{H_2O}$ . The latter are encapsulated in the former. Such assemblies we denote  $|\text{supra-}CD_{rot} < \text{supra-}CD_{elec}^{H_2O} >|$ . The state of  $H_2O$  belonging to both  $CD_{rot}$  and  $CD_{elec}^{H_2O}$  is a superposition of the state typifying the  $H_2O$  constituting  $CD_{rot}$  and the state typifying the  $H_2O$  constituting  $CD_{elec}^{H_2O}$ . In Figure 4 we present a schematic picture of bulk water containing macromolecules with sizable asymmetric charge distributions which stabilize  $|\text{supra-}CD_{rot} < \text{supra-}CD_{elec}^{H_2O} >|$ .



**Figure 4:** Schematic picture of bulk water containing macromolecules (represented by dark-blue balls) with a sizable asymmetric charge distributions at  $C < C_{crit}^{CD_{rot}}$ , i.e., at concentrations below the critical concentration for  $CD_{rot}$  formation. The water contains large elongated domains. These represent the  $10^{-5}$ - $10^{-4}$  m  $CD_{rot}$  stabilized by the macromolecules. The blue arrows represent the electric dipole moments of  $CD_{rot}$ , resulting from their ferroelectric ordered  $H_2O$ . Yellow-brown colored balls and their agglomerates represent, respectively,  $\sim 10^{-7}$  m sized  $CD_{elec}^{H_2O}$  and supra- $CD_{elec}^{H_2O}$ . These  $CD_{elec}^{H_2O}$  are located in  $CD_{rot}$ .  $CD_{elec}^{H_2O}$  are stabilized by  $CD_{rot}$ . The agglomerates of elongated domains and the yellow-brown balls represent the  $|\text{supra-}CD_{rot} < \text{supra-}CD_{elec}^{H_2O} >|$ . The various domains and macromolecules sizes are not presented according to their realistic scale ratios. The black cosine curves represent Far Infra Red (FIR) EMF facilitating interactions (symbolized by black double arrow) among  $H_2O$  over  $10^{-5}$ - $10^{-4}$  m distances. These interactions underlie  $CD_{rot}$  formation. The blue cosine curves represent UV EMF facilitating interactions (symbolized by blue double arrow) among  $H_2O$  over  $10^{-7}$  m distances. These interactions underlie  $CD_{elec}^{H_2O}$  formation.



**Properties of  $CD_{\text{plasma}}$  and  $IPD_{\text{plasma}}$**  --  $CD_{\text{plasma}}$  form when  $C_{\text{trans}}^{\text{IPDplasma}} < C < C_{\text{trans}}^{\text{CDplasma}}$  and  $IPD_{\text{plasma}}$  when  $C < C_{\text{trans}}^{\text{IPDplasma}}$ , due to interactions between solutes and tetra Herz to mega Herz EMF (Del Giudice *et al.*, 2000; Yinnon and Yinnon, 2012). At  $C_{\text{trans}}^{\text{IPDplasma}}$ ,  $CD_{\text{plasma}}$  transform into  $IPD_{\text{plasma}}$ . The transition concentrations  $C_{\text{trans}}^{\text{IPDplasma}}$  and  $C_{\text{trans}}^{\text{CDplasma}}$  depend on solute type and the polar solvent type. Typically  $\sim 10^{-6}$  M  $< C_{\text{trans}}^{\text{IPDplasma}} < \sim 10^{-4}$  M and  $C_{\text{trans}}^{\text{IPDplasma}} < C_{\text{trans}}^{\text{CDplasma}}$ .  $CD_{\text{plasma}}$  and  $IPD_{\text{plasma}}$  are composed of few solvated solutes and numerous polar solvent molecules -- these domains are not micelles.<sup>c</sup>  $CD_{\text{plasma}}$  and  $IPD_{\text{plasma}}$  are very stable domains. Energy gained by a solute on incorporation in  $IPD_{\text{plasma}}$  is larger than that of  $CD_{\text{plasma}}$  (few eV), implying this difference underlies solvation of difficult soluble compounds for  $C \leq C_{\text{trans}}^{\text{IPDplasma}}$ . The solvated solute molecules within  $IPD_{\text{plasma}}$  are crystalline-ordered, with the polar solvent molecules symmetrically aligned around their nearest neighbor solute molecule (see Figure 3b). *i.e.*, the *liquid*  $IPD_{\text{plasma}}$  have a crystalline structure.  $CD_{\text{plasma}}$ 's solvated solute molecules and their non-solvation shell solvent molecules locate randomly (see Figure 3a).

For a  $CD_{\text{plasma}}$ , the plasma oscillations of its identical solvated solute molecules are coherent, resonating in phase with the coherently condensed photons mediating their interactions (Del Giudice *et al.*, 2000). For an  $IPD_{\text{plasma}}$ , the plasma oscillations of its identical solvated solute molecules are in-phase, resonating in phase with the in-phase condensed photons mediating their interactions (Yinnon and Yinnon, 2012).

For the diameter of  $CD_{\text{plasma}}$  holds:  $\mathcal{D}^{\text{CDplasma}} \approx 1/\nu^{\text{CDplasma}}$ . For monovalent electro-

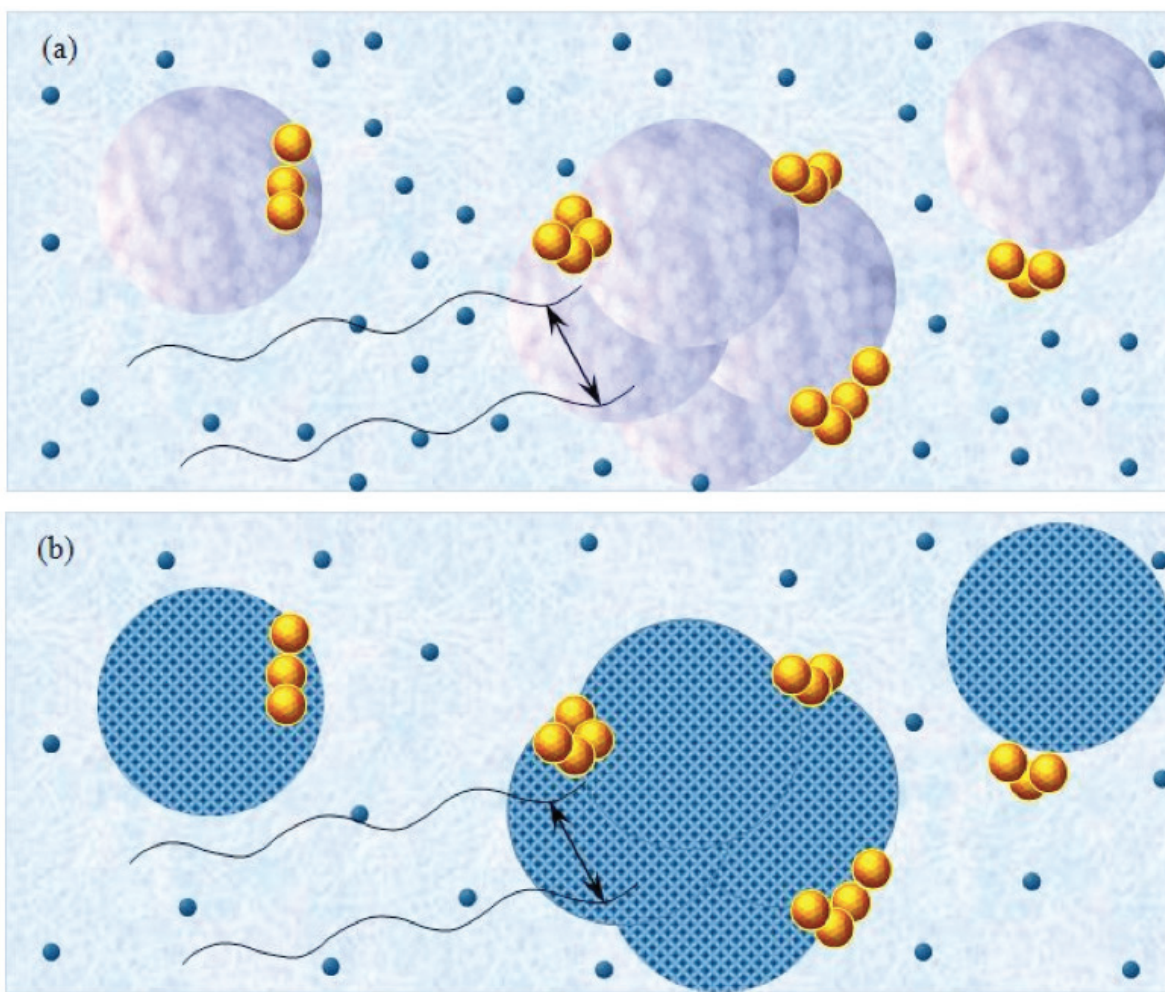
<sup>c</sup> A micelle is an aggregate of surfactant molecules. In aqueous solutions, its molecules' hydrophilic "head" regions are in contact with surrounding solvent, sequestering their hydrophobic single "tail" regions in the micelle's center.

lytes,  $\nu^{\text{CDplasma}} = 7.4 \times 10^3 (m_p/m_i)^{1/2} C^{3/4}$  GHz is the frequency of EFM mediating the attractive interactions between identical solvated solute molecules, with  $m_i$  and  $m_p$ , respectively, the mass of the ion and the proton. The expression of  $\nu^{\text{CDplasma}}$  results from  $C$  influencing a counter-ion mean charge density distribution  $\bar{\rho} \propto L_D^{-3/2}$ ,

$$\text{with: } L_D = \{[(\epsilon k_B T) / (8\pi e^2)]^{1/2} C^{-1/2}\}$$

denoting the Debye length,  $\epsilon$  is the dielectric constant, and  $e$  the charge of the electron.  $L_D$  equals the distance beyond which the Coulomb electric field around a solute molecule is at any instant fully screened by all its neighboring solvent molecules. For solutions containing non-monovalent ions or non-electrolytic solutes,  $\nu^{\text{CDplasma}}$  resembles the aforementioned expression, but is solute type dependent. The expressions for  $\nu^{\text{CDplasma}}$  connote that on dilution the diameter of  $CD_{\text{plasma}}$  increases, the number of  $CD_{\text{plasma}}$ 's solvated solute molecules diminishes and that of its solvent molecules enhances.

At  $C = C_{\text{trans}}^{\text{IPDplasma}}$ ,  $L_D$  equals the distance between closest neighbor's identical solute molecules. This distance is proportional to  $C^{-1/3}$ . For  $C > C_{\text{trans}}^{\text{IPDplasma}}$ ,  $L_D$  is smaller than this distance. For  $C \leq C_{\text{trans}}^{\text{IPDplasma}}$ ,  $L_D$  is larger than this distance. Thus for  $C \leq C_{\text{trans}}^{\text{IPDplasma}}$ , solvated solute molecules in  $CD_{\text{plasma}}$  start experiencing the repulsive Coulomb interactions with their identical nearest neighbors. Therefore on diluting below  $C_{\text{trans}}^{\text{IPDplasma}}$ , it is energetically unfavorable the number of solute molecules in  $IPD_{\text{plasma}}$  ( $N_s^{\text{IPDplasma}}$ ) and the frequency of their plasma oscillations decreases and the diameter of  $IPD_{\text{plasma}}$  increases. In other words at  $C \leq C_{\text{trans}}^{\text{IPDplasma}}$ : the frequency of  $IPD_{\text{plasma}}$ 's solutes' plasma oscillations and the diameter of  $IPD_{\text{plasma}}$  have values determined respectively, by  $\nu^{\text{CDplasma}}$  for  $C = C_{\text{trans}}^{\text{IPDplasma}}$  and  $N_s^{\text{IPDplasma}} = C_{\text{trans}}^{\text{IPDplasma}} \times V_{\text{trans}}^{\text{IPDplasma}}$  with  $V_{\text{trans}}^{\text{IPDplasma}}$  the volume of the domains at  $C_{\text{trans}}^{\text{IPDplasma}}$ . The aforementioned implies that *on diluting solutions below  $C_{\text{trans}}^{\text{IPDplasma}}$ , the number of  $IPD_{\text{plasma}}$  diminishes.*



**Figure 5a-b:** Schematic pictures of polar liquids containing solvated solutes. Figure (a) pertains to  $C^{IPD_{plasma}} < C < C^{CD_{plasma}}$  with purple-blue colored balls and their agglomerates representing, respectively,  $\sim 10^{-6}$  m  $CD_{plasma}$  and supra- $CD_{plasma}$ . Figure (b) pertains to  $C < C^{IPD_{plasma}}$ , with blue-crystalline structured balls and their agglomerates representing, respectively, the  $\sim 10^{-6}$  m crystalline structured  $IPD_{plasma}$  and supra- $IPD_{plasma}$ . Tiny dark-blue balls represent  $\sim 10^{-9}$  m solvated solutes with their hydration shells -- these randomly move in the interstices between the domains. Yellow-brown balls and their agglomerates represent, respectively,  $\sim 10^{-7}$  m  $CD_{elec}^{H_2O}$  and supra- $CD_{elec}^{H_2O}$ , which are stabilized by  $CD_{plasma}$  or  $IPD_{plasma}$ . Black cosine curves represent THz to MHz EMF facilitating interactions (symbolized by black double arrows) among the solvated solutes organized in  $CD_{plasma}$  or  $IPD_{plasma}$ . These interactions span  $\sim 10^{-6}$  m. Sizes of the domains and solvated solutes are not presented according to their realistic scale ratios.

$CD_{plasma}$  and  $IPD_{plasma}$  may stabilize  $CD_{elec}$  at their borders. Stabilization only occurs at the borders, because  $CD_{elec}$  cannot contain solutes. A schematic picture of a polar liquid with solutes at  $C_{trans}^{IPD_{plasma}} < C < C_{trans}^{CD_{plasma}}$ , i.e., wherein part of the solute molecules are organized in  $CD_{plasma}$ , is presented in Figure 5a. A similar picture but for  $C < C_{trans}^{IPD_{plasma}}$ , i.e., a solution wherein part of the solute mol-

ecules are organized in  $IPD_{plasma}$ , is presented in Figure 5b.

**Concentration ranges for CD formation** -- The abovementioned transition concentrations  $C_{crit}^{CD_{rot}}$ ,  $C_{trans}^{CD_{plasma}}$  and  $C_{trans}^{IPD_{plasma}}$  delineate  $C$  ranges for formation of, respectively,  $CD_{rot}$ ,  $CD_{plasma}$  and  $IPD_{plasma}$ . Their theoretical underpinnings vary greatly.

- Solvent molecules coherently oscillating between two of their rotational states is crucial for their organization in  $CD_{rot}$ . Solutes affect these oscillations. In the presence of few solutes, the rotational states of the solvent molecules get only slightly perturbed and the coherence persists. However, numerous solutes destroy the coherence (Del Giudice and Vitiello, 2006). Therefore  $CD_{rot}$  only form for  $C < C_{crit}^{CDrot}$ .

- Coherence of the plasma oscillations of solvated solute molecules sets in only below a solute type dependent  $C$ , *i.e.*,  $C_{trans}^{CDplasma}$  -- see Del Giudice *et al.* (2000) Eq. 14. The coherence is crucial for  $CD_{plasma}$  formation.

- In phase plasma oscillations of solvated solute molecules set in when nearest neighbor molecules experience their Coulomb interactions, *i.e.*, when their intermolecular distance is less than the Debye length -- a condition satisfied for  $C < C_{trans}^{IPDplasma}$  (Yinnon and Yinnon, 2012). These in phase oscillations are crucial for  $IPD_{plasma}$  formation.

- $H_2O$  coherently oscillating between their electronic ground state and an excited state is crucial for their organization in  $CD_{elec}^{H_2O}$ . Due to energetic reasons solutes cannot penetrate into  $CD_{elec}^{H_2O}$ . Thus formation of  $CD_{elec}^{H_2O}$  is independent of  $C$ .

**Superfluidic CD** --  $CD_{rot}$ ,  $IPD_{plasma}$  and  $CD_{elec}^{H_2O}$  are superfluidic domains (Bono *et al.*, 2012; Del Giudice *et al.*, 2013; Yinnon and Yinnon, 2012), *i.e.*, their molecules do not collide. A single collision would destroy their coherence.  $CD_{plasma}$  are not superfluidic -- their coherent oscillating solvated solute molecules and hydration shells' solvent molecules do not collide but their non-hydration solvent molecules do collide. Solvent molecules neither included in CD nor in solvated solutes' hydration shells move randomly and collide; so do solvated solute molecules not included in  $CD_{plasma}$  or  $IPD_{plasma}$ .

As to the impact of CD colliding with the

randomly moving molecules located in their interstices: when the collision energy is larger than the energy a molecule gains by its inclusion in CD, a molecule may desorb from the CD. The superfluidity implies that collisions excite the whole CD. For example, as in superfluidic Helium, whirlpools (so called rotons or vortexes) may be created. These excited CD states also may be induced by stirring or irradiation. When an excited state's energy is less than the energy a molecule gains by its inclusion in CD, it persists for macroscopic times. Hence CD constitute "long-term information storage devices". The superfluidity of CD has implications for the liquid's properties, *e.g.*, electric conductivity.

**Fractions of molecules included in CD** -- Fractions of solvent molecules included in  $CD_{rot}$ , in  $CD_{elec}$ , in  $IPD_{plasma}$  or  $CD_{plasma}$ , and fractions of solvated solutes included in  $CD_{plasma}$  or  $IPD_{plasma}$  depend on temperature, concentration and solute type (Preparata 1995 chapter 10; Del Giudice *et al.*, 2000; Yinnon and Yinnon, 2012).

## Discussion

Experimental data conforming to some predictions of the QED model for polar liquids and their solutions have been discussed in various publications. A detailed review of these is outside the scope of this paper, although indeed called for. Here we suffice with citing some relevant papers.

In their seminal publications predicting presence of  $CD_{rot}$ ,  $CD_{elec}$  and  $CD_{plasma}$  in aqueous systems, Del Giudice *et al.* (1988, 2000) and Arani *et al.* (1995) showed these clarify various phenomena which for many decades were unexplainable with the customary electrostatic models, *e.g.*, some of the so called anomalies of water. Yinnon and Yinnon (2009) point to significant experimental data in the literature pertaining to domains in aqueous solutions which have properties conforming to those of  $CD_{rot}$ ,



$CD_{elec}^{H_2O}$  and  $CD_{plasma}$ . Since 2009 additional experimental data corroborating the presence and properties of these domains has accumulated --- see for example Del Giudice *et al.*, 2010, 2013; De Ninno and Congiu Castellano, 2011; De Nino *et al.*, 2013; Germano *et al.*, 2013; Liu *et al.*, 2011, 2012; Yinnon and Yinnon, 2011; Yinnon and Elia, 2013; Elia *et al.*, 2015; Yinnon *et al.*, 2015c. Identification of  $IPD_{plasma}$  and their properties, together with analyses of some experimental data verifying these, were carried out by Yinnon and Yinnon (2012).

For all CD types, analyzed experimental data mainly concern their dynamics, their impact on the liquid's physicochemical properties and some of their structural aspects. As to the effects of EMF on CD, hitherto the related measurements have not been analyzed in depth, and the effects only have been inferred (Montagnier *et al.*, 2009; Yinnon and Yinnon, 2011; Yinnon and Elia, 2013).

## Conclusions

This paper relates on the structure of polar liquids and their solutions. The quantum electrodynamics (QED) theory's predictions for the structure of these liquids considerably differ from those of the customary electrostatic theories. While the former explicitly describe the electrodynamic interactions mediated by EMF, the latter assume that these interactions can be treated as small perturbations or ignored. Since the former shows that electrodynamic interactions under certain conditions induce phase transitions leading to formation of coherence domains, the limits of the customary theories have to be carefully assessed. Hitherto, only few assessments were undertaken. These mainly focus on the domains' dissipative dynamics and their impact on the liquid's physicochemical properties.

A detailed review of these assessments is called for - in particular it is needed to map

the many challenging research projects still required for verifying numerous aspects of polar liquids predicted by QED. However, even now it is overt that QED analyses are required for a series of experimental data published since 2008 pertaining to direct observation of domains in polar liquids and the effects of EMF on these, *e.g.*, for liquids perturbed by serial dilutions or immersion of membranes [for reviews of the data see Konovalov and Ryzhkina (2014b) and Elia *et al.* (2015)]. As to the liquids perturbed by membranes, some QED analyses are presented in Del Giudice *et al.* (2010, 2013) and Yinnon *et al.* (2015c). As to serial diluted solutions, analyses of their directly observed domains and the impact of EMF, these are the foci of the following papers in this journal's volume (Yinnon and Liu, 2015a and 2015b).

Though the aforementioned stresses the need for gathering additional experimental data supporting the predictions of the QED model, it should be emphasized that additional theoretical work is required too. Hitherto only a few properties of polar liquids resulting from their electrodynamic interactions have been derived. For example, a mathematical expression for the critical concentration below which  $CD_{rot}$  form is still lacking. Also the derivation of quantitative expressions for physicochemical variables of polar liquids containing CD, *e.g.*, their electric conductivity, dielectric permittivity and surface tension, are important challenges ahead. Detailed predictions on the excited states of the various CD and their influence on spectra of polar liquids also has to be addressed in future research projects.

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**Table 1:** List of abbreviations in alphabetic order, followed by Greek symbols abbreviations.

Abbreviations	Explanation
$C$	Concentration
$C_{\text{crit}}$	Critical concentration
$C_{\text{crit}}^{\text{CDrot}}$	Critical concentration below which $\text{CD}_{\text{rot}}$ may form
$C_{\text{trans}}^{\text{CDplasma}}$	Transition concentration for $\text{CD}_{\text{plasma}}$ formation
$C_{\text{trans}}^{\text{IPDplasma}}$	Transition concentration for $\text{IPD}_{\text{plasma}}$ formation
$C_{\text{thr}}$	Threshold concentration
$\text{CD}$	Coherence domain
$\text{CD}_{\text{elec}}^{\text{H}_2\text{O}}$	Coherence domains composed of water molecules coherently transiting between their electronic ground state and an excited state
$\text{CD}_{\text{plasma}}$	Coherence domain composed of few solvated solutes performing coherent plasma oscillation and numerous polar solvent molecules
$\text{CD}_{\text{rot}}$	Coherence domains of ferroelectric ordered polar solvent molecules
$\mathcal{D}^{\text{CDplasma}}$	Diameter of $\text{CD}_{\text{plasma}}$
EMF	Electro-magnetic fields
eV	Electron Volt
$\text{H}_2\text{O}$	Water molecule
$\text{IPD}_{\text{plasma}}$	In phase domains composed of few solvated solutes and numerous polar solvent molecules performing in phase plasma oscillation.
K	Kelvin
$k_B$	Boltzmann constant
$L_D$	Debye length
M	Molarity in mol per liter
m	meter
$N$	Number of molecules
$N_s^{\text{IPDplasma}}$	Number of solute molecules in $\text{IPD}_{\text{plasma}}$
$P_r$	Probability
$P_N$	Probability photon excites molecule
QED	Quantum electro-dynamics
SDVSPL	Serial diluted vigorous shaken polar liquids
Supra-CD	Agglomerate of coherence domains
Supra- $\text{CD}_{\text{rot}}$ <supra- $\text{CD}_{\text{elec}}^{\text{H}_2\text{O}}$ >	Agglomerate of $\text{CD}_{\text{rot}}$ containing agglomerates of $\text{CD}_{\text{elec}}^{\text{H}_2\text{O}}$
$T$	Temperature in degree Kelvin
ULC	Ultra low concentration
$V_e$	Volume of ensemble of molecules
$V_{\text{photon}}$	Volume spanned by photon with wavelength $\lambda$
$V_{\text{trans}}^{\text{IPDplasma}}$	Volume of an $\text{IPD}_{\text{plasma}}$ at $C_{\text{trans}}^{\text{IPDplasma}}$
$\lambda$	Wavelength of photon
$\nu^{\text{CDplasma}}$	Frequency of photon trapped in $\text{CD}_{\text{plasma}}$

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